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Lecture for M. Sc. Physics II Semester students

Paper-IV: Electronic Devices

Unit-1 Transistors and Microwave Devices

Transfer of electrons from Central Valley to the satellite valley is called Transfer Electron Mechanism. (5)

Gunn Diodes (Transfer Electron Devices)

With the help of junction diodes high microwave power can not be obtained.

While studying the properties of thin specimens of GaAs, J. B. Gunn in 1963 discovered that under high electric stress there is periodic fluctuations in the current passed by the material. This effect is known as Gunn effect or as bulk-transferred electron effect. T.E.D are bulk-effect devices that utilizes hot electrons to produce a voltage control differential negative resistance.

The phenomena that gives rise to differential negative resistance is the decrease in electron mobility with the electric field.

This decrease is caused by the transfer of conduction electron from high mobility low energy state to a low mobility high energy state under the action of high electric fields.

Fig (A) shows the construction of a typical Gunn diode when a small D.C voltage is applied across thin slices of GaAs, negative resistance manifests itself under certain conditions.

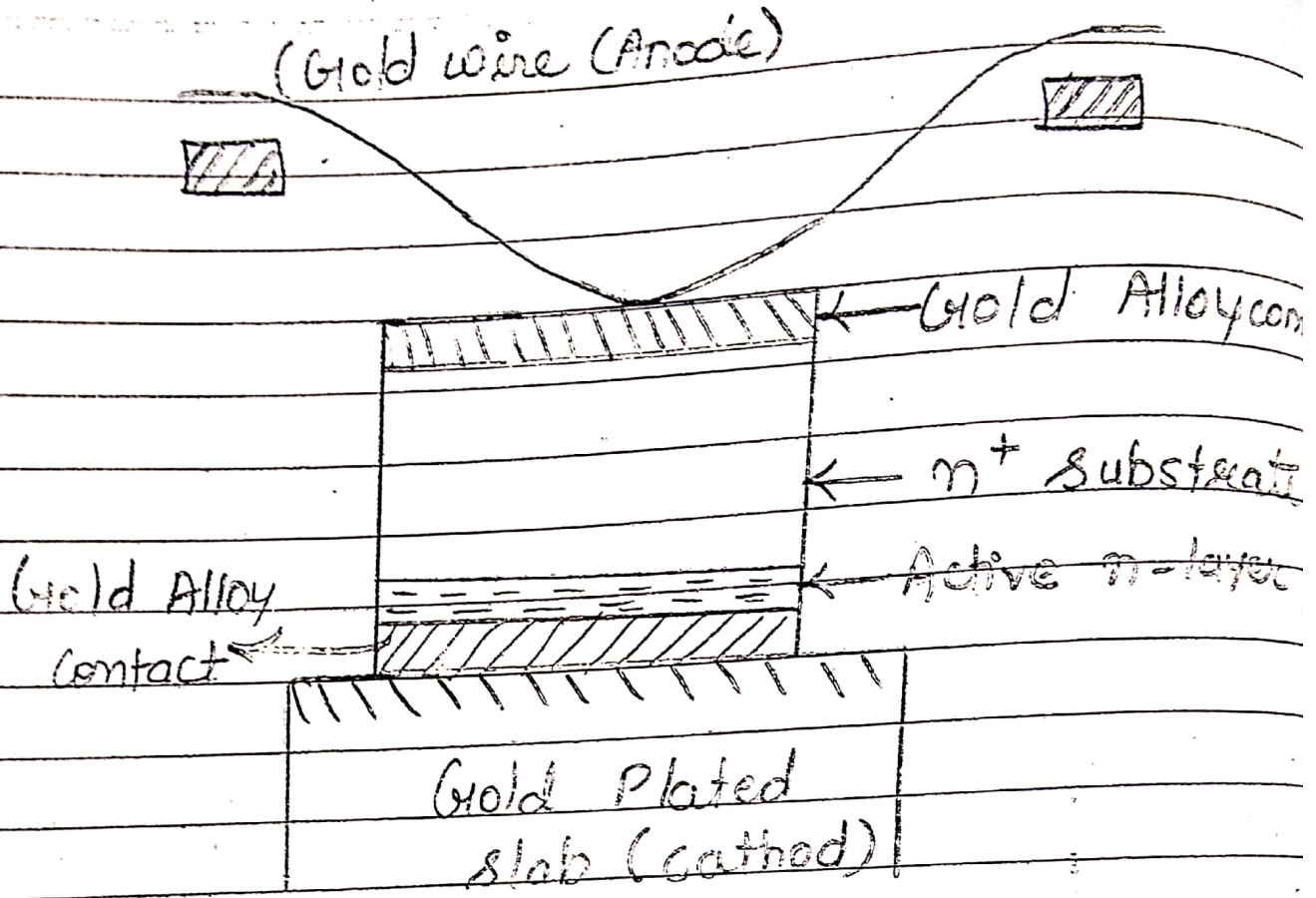


Fig. (A)

If a slice is ~~connected~~ to the suitably tuned circuit oscillation occurs. Hence the slice is very thin, voltage gradient across the slice become very high. Due to this reason the electron velocity is also very high which causes the oscillation to occur at microwave frequency.

Fig (b) shows the cross-section to the E-k diagram of the GaAs conduction band. The diagram has following features -

- 1) The lowest minimum occurs at $k=0$ (valley 1) here, the E-k diagram has

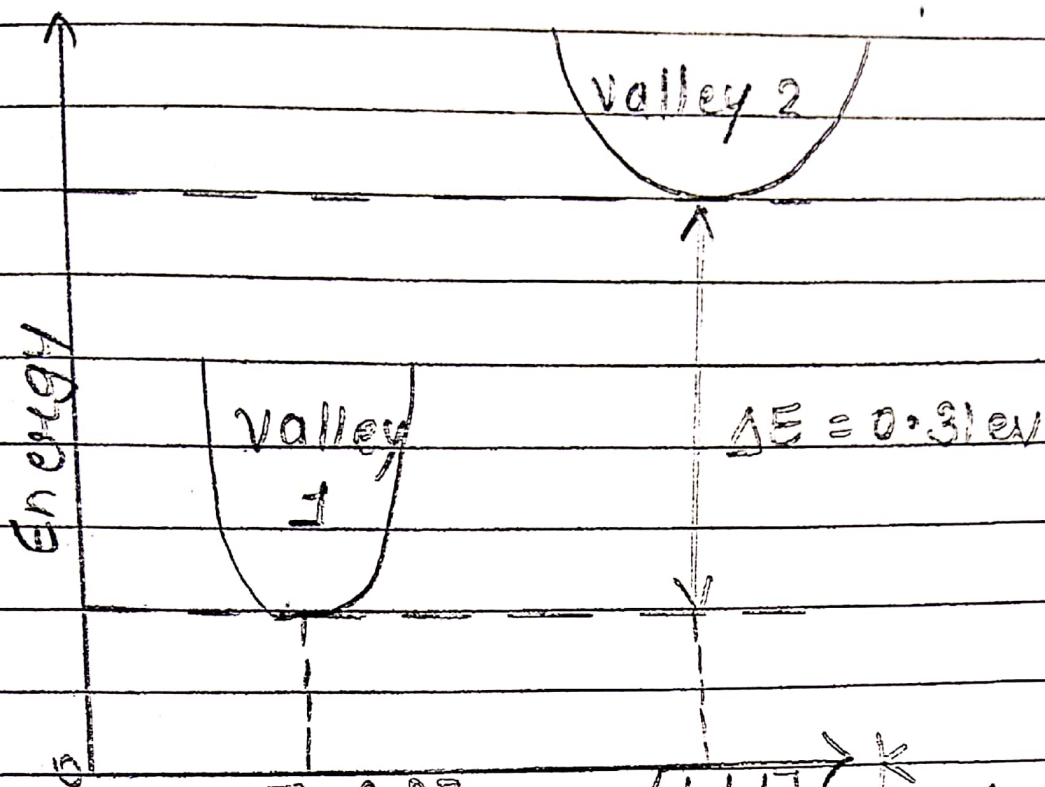


Fig (b) E-K diagram of GaAs conduction band

a sharp curvature, and the electrons in this lower valley has a low effective mass m_1 , and hence high mobility μ_1 . Besides the lower valley there are energy maxima along $[1,1,1]$ axis known as satellite valley each of which is separated from the lower valley by an energy $\Delta E = 0.31 \text{ eV}$. In a satellite valley the electron effective mass m_2 is high and the mobility μ_2 considerably lower than μ_1 .

2) The density of state in satellite valley is considerably higher than that in

the lower valley.

3) The energy difference ΔE is large compared to the thermal energy kT of the e^- at the room temperature. Thus the transfer of e^- from the lower to the satellite valley by thermal agitation is not very likely. Furthermore, since ΔE is small compared to energy gap E_g , electron transfer can occur at field much lower than those required for avalanche breakdown.

Consider now that an electric field \vec{E} is applied to a sample of GaAs. At low fields all the conduction electrons are located in the lower valley. However, as \vec{E} is increased, electrons gain energy from the field and make a transition to the low mobility satellite valleys. As a result, the conductivity of the material decreases at high fields. The electron transfer sets in abruptly after the field has reached the threshold value E_T . This causes a decrease in the average drift velocity of the electron and results in a region of negative

differential mobility (NDM) - The conductivity in a two-valley semiconductor can be written as

$$\sigma = q [n_1(\vec{E}) \mu_1 + n_2(\vec{E}) \mu_2] = q n_0 \frac{v_n(\vec{E})}{E} \quad (1)$$

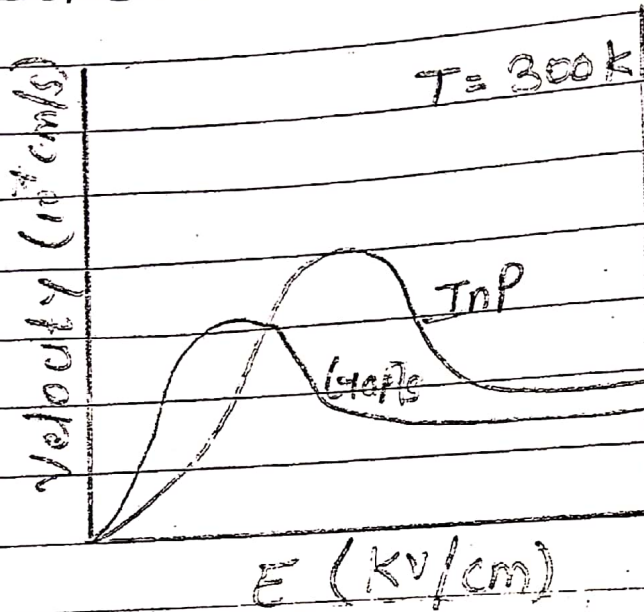
where n_1 and n_2 are the electron concentrations in the lower and the upper valley respectively and $n_0 = (n_1 + n_2)$. From eqn (1), the average drift velocity $v_n(\vec{E})$ is obtained as

$$v_n(\vec{E}) = \frac{[n_1(\vec{E}) \mu_1 + n_2(\vec{E}) \mu_2] E}{n_0} \quad (2)$$

It can be seen from this relation that the differential mobility $\frac{dv_n}{dE}$ can become negative at a sufficiently high field. Since n_1 decreases and n_2 increases with E . At very high fields when all the electrons have made a transition to the upper valley, $v_n(\vec{E})$ again starts increasing, with E making $\frac{dv_n}{dE}$ positive.

Static I-V characteristic of a DNM crystal

Let us consider a sample of length l of a semiconductor whose velocity field characteristic is shown in fig.



Let voltage V be applied to the sample. If V is increased gradually, we can expect to obtain the I-V characteristic similar to that of a tunnel diode, provided the field is uniform throughout. However in real crystals, local fluctuations in the mobile carrier concentration occur due to doping inhomogeneity. This causes the field to become non-uniform and as a result, the negative resistance characteristic is not observed. The one-dimensional equations governing the

I-V characteristic of the sample obeys the Poisson equation.

$$\frac{\partial E}{\partial x} = \frac{q}{\epsilon_s} [n(x) - n_0] \quad \text{--- (3)}$$

and the current density equation

$$J(x) = q n(x) v_n(x) \quad \text{--- (4)}$$

Where, n_0 = thermal equilibrium electron concentration

$n(x)$ = concentration at any point x
 From the above two relations we obtain

$$n(x) = \frac{J(x)}{v_n(x)q}$$

∴ eqⁿ (3) becomes

$$\frac{\partial E(x)}{\partial x} = \frac{q}{\epsilon_s} \left[\frac{J(x)}{v_n(x)q} - n_0 \right] \quad \text{--- (5)}$$

A computer solution of this non-linear differential equation reveals that in the steady state, $E(x)$ increases monotonically with x and no negative resistance characteristic is observed. This is because an internal space charge exists in the sample and $n(x)$ varies with E such