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Classical mechanics.

Uniformly Rotating Frames: →

We know that the earth itself rotates about its axis in 24 hours. Therefore, any frame fixed with the earth will also rotate with it and so it will be a noninertial frame.

Suppose that a frame $s'(x_r, y_r, z_r)$ is rotating with an angular velocity ω relative to an inertial frame $s(x_i, y_i, z_i)$. For simplicity, we assume that both of the frames have common origin O and common z -axis. In case of the earth, the common origin O may be considered as the centre of the earth, z -axes as coinciding with its rotational axis and the frame s' as rotating with earth relative to the non rotating frame s .

The position vector of a particle P in both frames will be the same, i.e. $R_i^o = R_r^o = R$, because the origins are

Coincident now, if the particle p is stationary in the frame s , the observer in the rotating frame s' will see that the particle is moving oppositely with linear velocity $-\omega \times R$. Thus, if the velocity of the particle in the frame s is $\left(\frac{dR}{dt}\right)_i^o$, then the velocity $\left(\frac{dR}{dt}\right)_{s'}$ in the rotating frame will be given by

$$\left(\frac{dR}{dt}\right)_{s'} = \left(\frac{dR}{dt}\right)_i^o - \omega \times R$$

or $\left(\frac{dR}{dt}\right)_i^o = \left(\frac{dR}{dt}\right)_{s'} + \omega \times R \quad \text{--- (1)}$

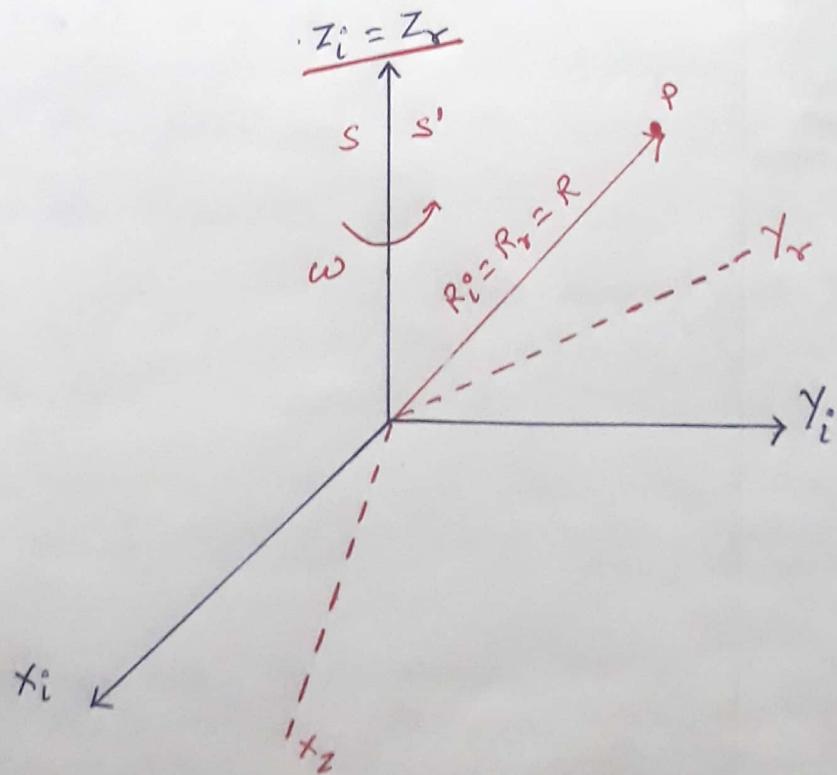


Figure: Uniformly Rotating Frame

In fact this equation holds for all vectors and relates the time derivatives of a vector in the frame S and S'. Therefore, relation (1) may be written in the form of operator eqn

$$\left(\frac{d}{dt}\right)_i = \left(\frac{d}{dt}\right)_r + \omega \times R \quad \text{--- (2)}$$

Writing $\frac{dR}{dt} = v$ for the velocity of the particle, we have

$$v_i^o = v_r + \omega \times R \quad \text{--- (3)}$$

Now, if we operate eqn (2) on velocity vector v_i^o , we have

$$\left(\frac{d v_i^o}{dt}\right)_i = \left(\frac{d v_r}{dt}\right)_r + \omega \times v_i^o$$

Substituting the value of v_i^o in the right hand side of this relation from eqn (3), we obtain

$$\begin{aligned} \left(\frac{d v_i^o}{dt}\right)_i &= \left[\frac{d}{dt} (v_r + \omega \times R) \right]_r + \omega \times (v_r + \omega \times R) \\ &= \left(\frac{d v_r}{dt}\right)_r + \frac{d\omega}{dt} \times R + \omega \times \left(\frac{dR}{dt}\right)_r + \omega \times v_r + \omega \times (\omega \times R) \end{aligned}$$

If we write the acceleration

$$\frac{dv}{dt} = a \quad \text{and} \quad \left(\frac{dR}{dt}\right)_r = v_r, \text{ then}$$

$$\alpha_i = \alpha_r + 2\omega \times v_r + \omega \times (\omega \times R) + \frac{d\omega}{dt} \times R$$

For earth, ω is constant so $\frac{d\omega}{dt} = 0$, then

$$\alpha_i = \alpha_r + 2\omega \times v_r + \omega \times (\omega \times R) \quad (4)$$

If m is the mass of the particle then the force in the rotating frame is

$$m\alpha_r = m\alpha_i - 2m\omega \times v_r - m\omega \times (\omega \times R)$$

But $m\alpha_r = F_p + F_0$, therefore fictitious force F_0 is given by

$$F_0 = -2m\omega \times v_r - m\omega \times (\omega \times R) \quad (5)$$

where $-2m\omega \times v_r$ is the Coriolis force and $-m\omega \times (\omega \times R)$, the Centrifugal force

The centrifugal force is the only fictitious force, acting on a particle which is rest ($v_r = 0$) in the rotating frame. The Centrifugal force can be written as

$$-m\omega \times (\omega \times R) = m\omega^2 r \quad (6)$$

The Coriolis and Centrifugal forces are often called fictitious forces because the bodies don't actually feel them, they are due only to the observer's motion.