M.Sc. II Sem

**PHY-201** [Atomic and Molecular Physics]

### Unit- III Nuclear Magnetic Spectroscopy



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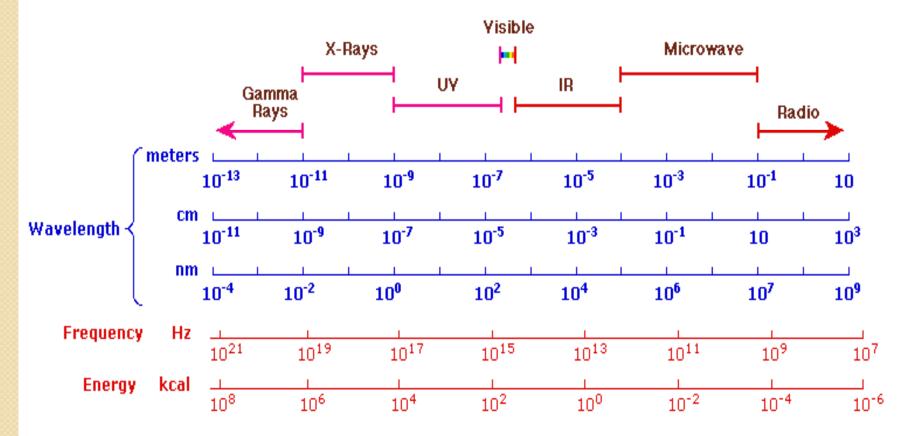
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Outline

### Introduction to Spectroscopy Properties of Spin The Nuclear Magnetic Moment The Nuclear Spin Quantum Number Magnetic Properties of atomic Nuclei Nuclear precession in Magnetic field Semi-Classical Description Quantum mechanical Description Quantum- classical correspondence

## Introduction to Spectroscopy

Spectroscopy is the study of the interaction of electromagnetic radiation (light) with matter.



The Electromagnetic Spectrum

NMR uses electromagnetic radiation in the radio frequency range

- •Long wavelength, very low energy
- •Low energy has significant consequences:
- •Sharp signals (Good)
- •Poor sensitivity (Bad)
- Longer experiment time (Bad)



# **Properties of Spin**

Spin is a fundamental property of nature

- Any unpaired electron, proton, or neutron will possess a spin of  $\frac{1}{2}$
- Atomic nuclei, which are composed of protons and neutrons, may also possess spin
- The spin of an atomic nucleus is determined by the number of protons and neutrons
- Atoms with odd number of protons will have spin
- Atoms with odd number of neutrons will have spin
- Atoms with EVEN number of protons and neutrons will not have spin

The value of the nuclear spin is defined by **I**, the nuclear spin quantum number and can have values of (I = 0, 1/2, I, 3/2, 2, 5/2, ...)

A nucleus of spin I can exist in (2I+I) spin states.

### The Nuclear Magnetic Moment

•All atomic nuclei can be characterized by a nuclear spin quantum number, *I. I can be*  $\geq 0$  and any multiple of  $\frac{1}{2}$ 

•Nuclei with I= 0 do not possess nuclear spin and consequently are termed 'NMR silent'.

•All nuclei with I≠0 possess spin, charge and angular momentum P, resulting in a nuclear magnetic moment μ.

$$\mu = \gamma P$$

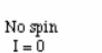
Where  $\gamma$  is the gyromagnetic atio of the nucleus.

Values of  $\gamma$  can be positive or negative and determine the sense of precession and thus the direction of the magnetic moment.

### I = the nuclear spin quantum number







Spinning sphere I = 1/2



I = 1, 3/2, 2, ....

For Nuclei of:

Example

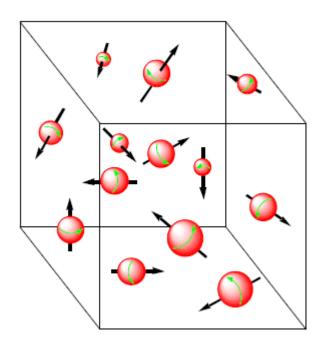
Odd Mass	Half Integer	<sup>1</sup> H, <sup>13</sup> C
Even Mass/Even Charge	Zero	$^{12}C, ^{16}O$
Even Mass/Odd Charge	Integer	${}^{2}\text{H}, {}^{14}\text{N}$

If I = 0, NMR Inactive If  $I \ge 1$ , Quadrupolar (non-spherical nuclear charge distribution)

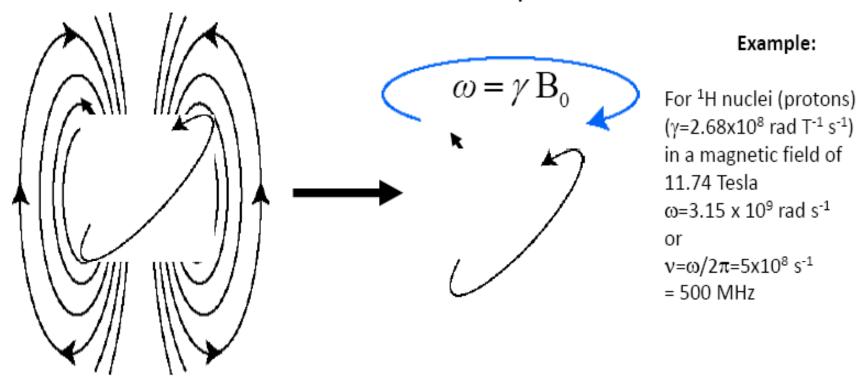


### Magnetic properties of atomic nuclei

- The magnetic moment (µ)is a vector quantity that has both magnitude and direction
- In the absence of an external magnetic field the magnetic moments (µ)are randomly orientated.

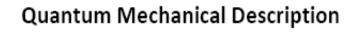


Semi-Classical Description



Magnetic Field B<sub>0</sub>

The Magnetic Field ( $\mathbf{B}_0$ ) exerts torque on angular momentum ( $\mathbf{L}$ ) and causes Nuclear Precession, analogous to precession of spinning top. The frequency of the precession ( $\boldsymbol{\omega}$ ), often called the Larmor frequency, is proportional to the gyromagnetic ration ( $\gamma$ ) and the strength of the external magnetic field ( $\mathbf{B}_0$ ).

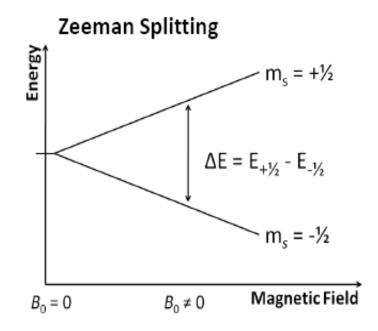


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Во

In the macroscopic world the two magnets can be aligned in an infinite number of orientations. At the atomic level, these alignments are **quantized** and the number of orientations (spin states) are equal to **21+1**. We will only deal with spin ½ nuclei.

The different **quantized** orientations will each have an energy level determined by the Zeeman splitting

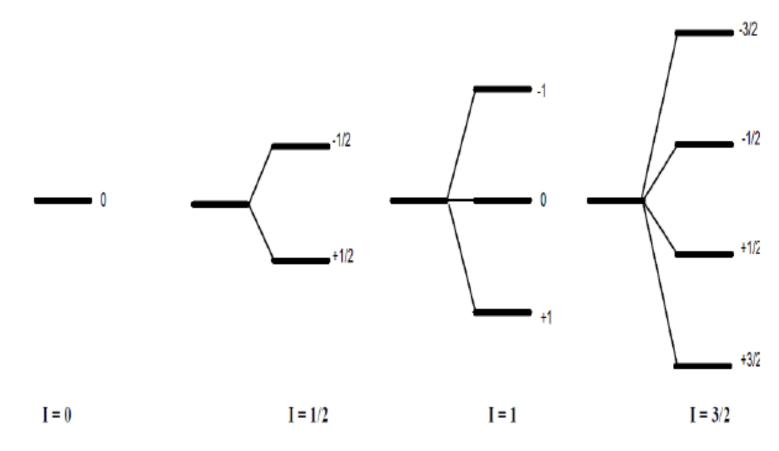


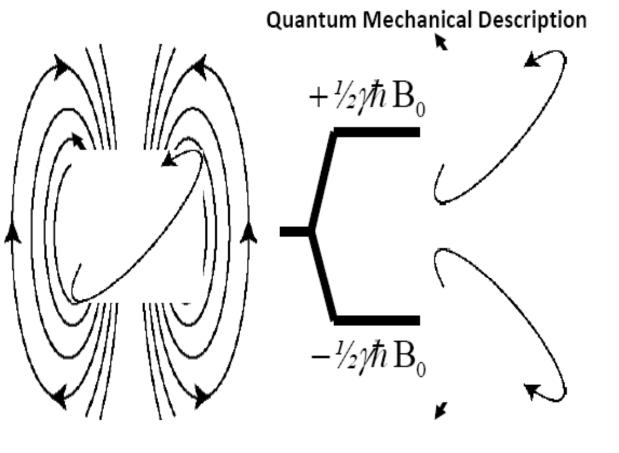


**Quantum Mechanical Description** 

The energy levels are more complicated for I > 1/2

Zeeman Splitting





 $\Delta E = \gamma \hbar B_0$ 

Knowing ∆E, we can stimulate the transition between these two states by applying an RF field such that:

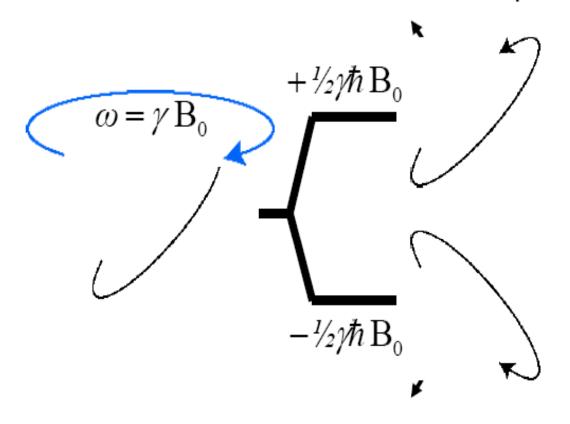
 $E = hv = \hbar\omega$ 

Magnetic Field B<sub>0</sub>

Zeeman Splitting

In NMR spectroscopy we are going to perturb the spin states by stimulating transitions between the energy levels.

**Quantum-Classical Correspondence** 



$$\Delta E = \gamma \hbar B_0$$

Knowing ∆E, we can stimulate the transition between these two states by applying an RF field which satisfies the resonance condition:

$$E = \hbar \omega = \gamma \hbar B_0$$

# Thank You