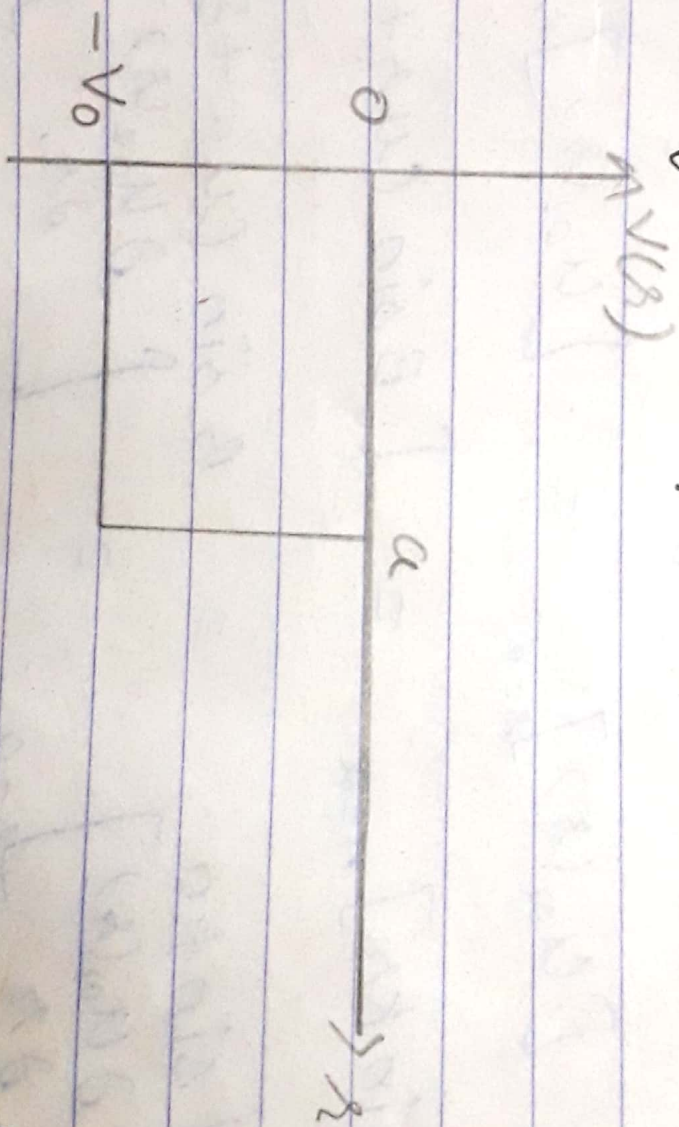


Scattering By a Square Well Potential -



(9)

We consider the S-wave scattering by a square well

$$V(r) = -V_0 \quad \text{for } 0 < r < a$$

$$V(r) = 0 \quad \text{for } r > a$$

For  $l=0$ , the radial part of Schrodinger eq. would be

Case (1)  $\rightarrow \frac{d^2 u_{01}}{dr^2} + k^2 u_{01}(r) = 0 \quad \text{--- (1)}$   
 for  $0 < r < a$

Where  $k^2 = \frac{2m}{\hbar^2} (E + V_0)$  --- (2)

Case (2)  $\rightarrow \frac{d^2 u_{02}}{dr^2} + \mu^2 u_{02}(r) = 0 \quad \text{--- (3)}$   
 for  $r > a$

Where  $\mu^2 = \frac{2mE}{\hbar^2}$

The general solu. of eq. (1) is

$$u_{01}(r) = A \sin kr + B \cos kr \quad \text{--- (4)}$$

The cosine solu. will make  $R_0(r) = u_{01}(r)/r$  go to  $\infty$  at  $r=0$

only acceptable solu. is

$$u_{01}(r) = A \sin kr \quad \text{--- (5)}$$

Similarly the solu. of eq. (3) can be taken as

$$u_{02}(r) = B \sin(\mu r + \delta_0) \quad \text{--- (6)}$$

Because the difference will be only in phase factor. Putting the boundary condit<sup>n</sup> at  $r=a$ .

$$[u_{01}(r)]_{r=a} = [u_{02}(r)]_{r=a}$$

$$[A \sin kr]_{r=a} = [B \sin(\mu r + \delta_0)]_{r=a}$$

$$A \sin ka = B \sin(\mu a + \delta_0) \quad \text{--- (7)}$$

$$\& \left[ \frac{\partial u_{01}(r)}{\partial r} \right]_{r=a} = \left[ \frac{\partial u_{02}(r)}{\partial r} \right]_{r=a}$$

(11)

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$$\left[ \frac{\partial}{\partial r} A \sin kr \right]_{r=a} = \left[ \frac{\partial}{\partial r} B \sin(\mu r + \delta_0) \right]_{r=a}$$

$$[A \cdot k \cos kr]_{r=a} = [B \cdot \mu \cos(\mu r + \delta_0)]_{r=a}$$

$$A k \cos ka = B \mu \cos(\mu a + \delta_0) \quad \text{--- (8)}$$

Dividing eq. (7) by (8)

$$\frac{\sin ka}{k \cos ka} = \frac{\sin(\mu a + \delta_0)}{\mu \cos(\mu a + \delta_0)}$$

$$\Rightarrow \tan(\mu a + \delta_0) = \frac{\mu}{k} \tan ka$$

$$\mu a + \delta_0 = \tan^{-1} \left\{ \frac{\mu}{k} \tan ka \right\}$$

$$\delta_0 = \tan^{-1} \left\{ \frac{\mu}{k} \tan ka \right\} - \mu a \quad \text{--- (9)}$$

$$\therefore \sigma = \frac{4\pi}{k^2} \sin^2 \delta_0$$

$$\sigma = \frac{4\pi}{k^2} \left[ 1 + \frac{1}{\tan^2 \delta_0} \right]^{-1} \quad \text{--- (10)}$$

$$\mu a \ll 1$$

In eq. (9) we notice that  $\tan \delta_0$  would vanish if

$$\boxed{\frac{\tan ka}{k} = \frac{\tan \mu a}{\mu}}$$

We are assuming  $\mu a \ll 1$  so that  $\delta_1, \delta_2$  etc. are negligible.  
This known as RAMSAUER effect.