

**APPLICATION  
OF  
SCHRODINGER EQUATION  
TO UNDERSTAND THE HYDROGEN ATOM**

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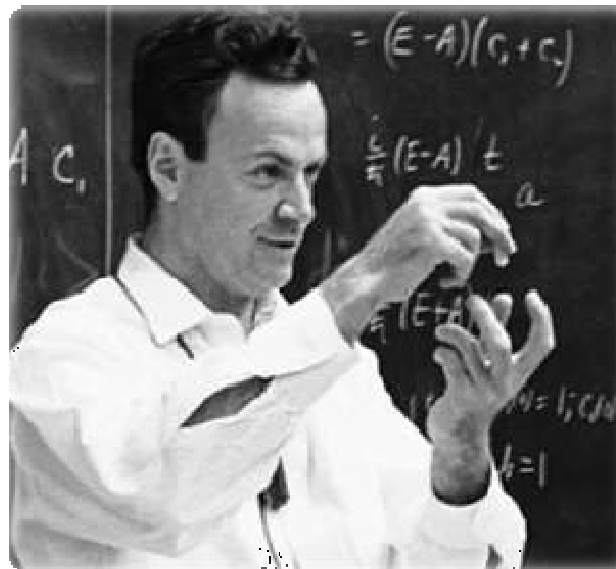
**SCHOOL OF STUDIES IN CHEMISTRY  
& BIOCHEMISTRY**



**“I think I can safely say that nobody understands quantum mechanics.”**

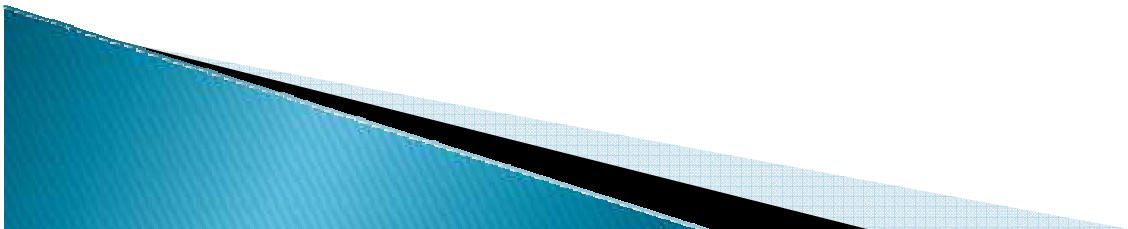
Ch. 6, “Probability and Uncertainty”

–Richard Feynman



# CONTENTS

- Wave function and its physical significance.
- Schrödinger time independent equation.
- Application of the Schrödinger Equation to the Hydrogen Atom.
- Solution of the Schrödinger Equation to the Hydrogen Atom.
- Solutions to radial, angular and azimuthal equation.
- Quantum Numbers.
- Normalized Spherical Harmonics

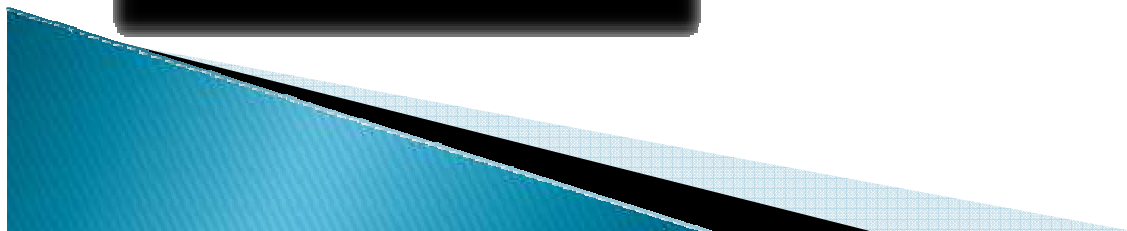
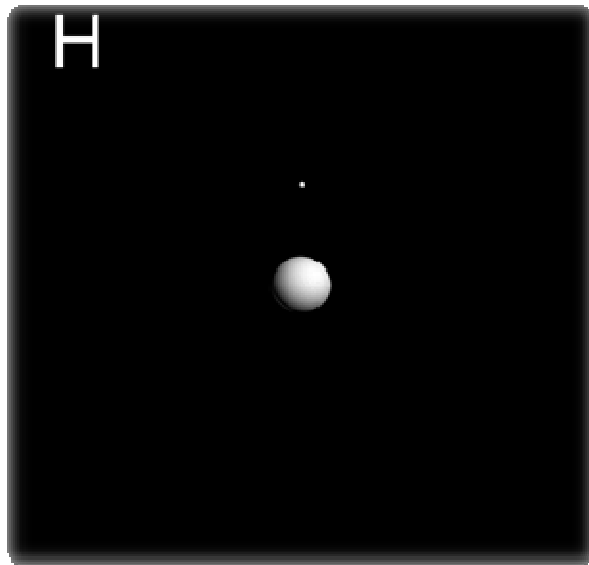


**The whole aim of this project is to convert ...**

**This**

**Into**

**This**



- **WAVE FUNCTION AND IT'S PHYSICAL SIGNIFICANCE (  $\psi$  )**



The physical significance or interpretation of  $|\psi|^2$  or  $\psi \psi^*$  was given by Max Born in 1927.

1.  $\psi^2$  or  $\psi \psi^*$  represents probability per unit volume, of finding a particle described by a wave function at a particular point (x,y,z) at time 't' in that volume.

i.e. 
$$\psi^2 = \frac{p}{dV}$$

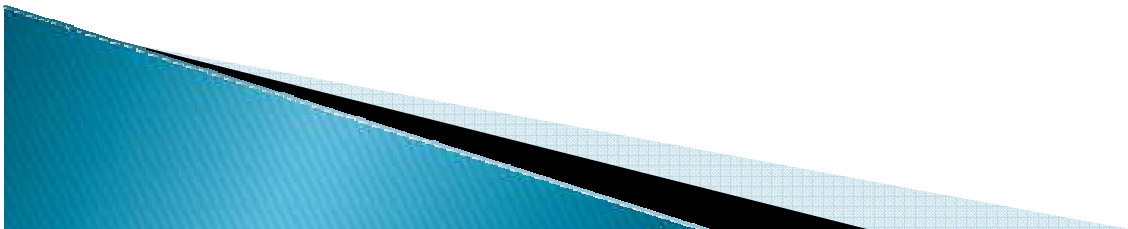
Max Born

# IMPORTANT CHARACTERISTICS OF A WAVE FUNCTION ASSOCIATED WITH WAVE

1. Wave function should be normalised. i.e. it should satisfy the following condition:

$$\int_{-\infty}^{\infty} \psi^2 dx dy dz = 1$$

2. Wave function should be finite.  $\infty$  probability has no meaning.
3. It should have a single value. (as at particular point probability can have only one value.)



# SCHRODINGER TIME INDEPENDENT EQUATION

Schrodinger used analogy of stationary waves on string for matter waves.

For a particle having position coordinates (x,y,z) and having wave function  $\psi$  equation of matter wave is

$$\frac{d^2\psi}{dt^2} = u^2 \nabla^2 \psi \quad , \quad \text{where } \nabla^2 = \left( \frac{d^2}{dx^2} + \frac{d^2}{dy^2} + \frac{d^2}{dz^2} \right) \\ = \text{Laplacian Operator}$$

The solution of this equation is given as,

$$\psi = \psi_0 e^{-i\omega t} \quad , \quad \text{where } \omega = 2\pi\nu \text{ and } u = \nu\lambda$$

Double differentiating and substituting in above equation we get,

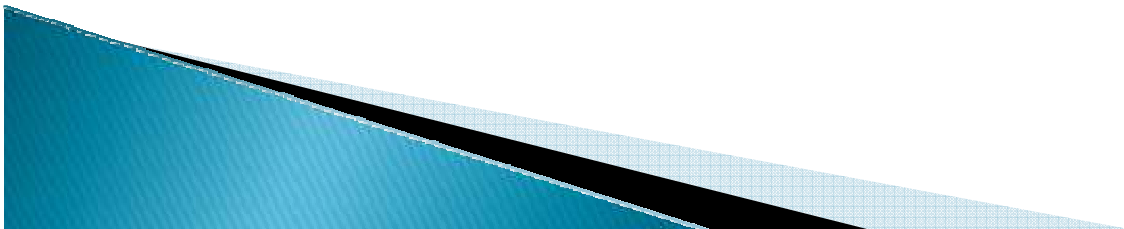
$$\nabla^2 \psi + \frac{4\pi^2}{\lambda^2} \psi = 0$$

On further simplification we get,

$$\nabla^2 \psi + \frac{8\pi^2 m}{h^2} (E - V) \psi = 0$$

Where,

- h = Planck's constant
- m = mass of particle
- E = total energy of particle
- V = potential energy of particle





# Application of the Schrödinger Equation to the Hydrogen Atom

- The approximation of the potential energy of the electron-proton system is electrostatic:

$$V(r) = -\frac{e^2}{4\pi\epsilon_0 r}$$

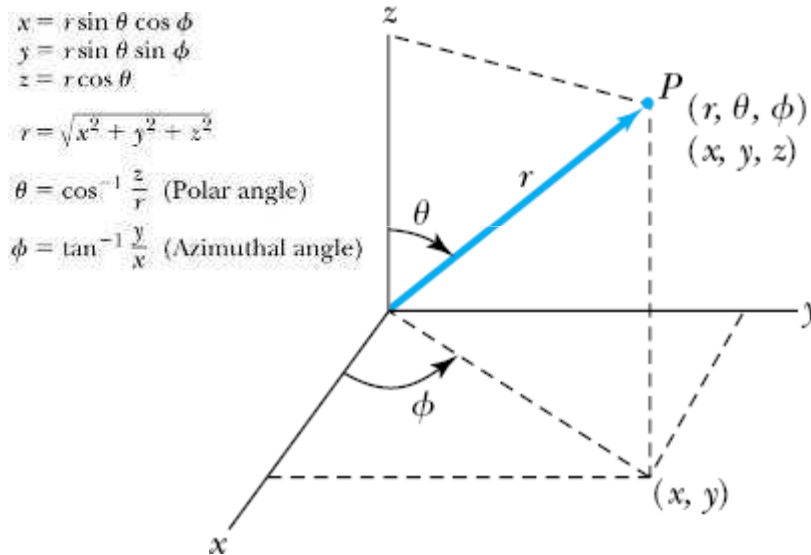
- Rewrite the three-dimensional time-independent Schrödinger Equation.

$$-\frac{\hbar^2}{2m} \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \psi(\vec{r}) + \left( -k \frac{Ze^2}{r} \right) \psi(\vec{r}) = E\psi(\vec{r})$$

- ▶ For Hydrogen-like atoms (He<sup>+</sup> or Li<sup>++</sup>)
- Replace  $e^2$  with  $Ze^2$  ( $Z$  is the atomic number).
- Use appropriate reduced mass  $\mu$ .

# Application of the Schrödinger Equation

- ▶ The potential (central force)  $V(r)$  depends on the distance  $r$  between the proton and electron.



Transform to spherical polar coordinates because of the radial symmetry.

Insert the Coulomb potential into the transformed Schrödinger equation.

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} + \frac{2\mu}{\hbar^2} (E - V) \psi = 0$$

## Application of the Schrödinger Equation

- ▶ The wave function  $\psi$  is a function of  $r, \theta, \phi$  .
  1. Equation is separable.
  2. Solution may be a product of three functions.

$$\psi(r, \theta, \phi) = R(r)f(\theta)g(\phi)$$

- ▶ We can separate Equation into three separate differential equations, each depending on one coordinate:  $r, \theta$ , or  $\phi$  .

# Solution of the Schrödinger Equation for Hydrogen

- ▶ Substitute into Equation and separate the resulting equation into three equations:  $R(r)$ ,  $f(\theta)$ , and  $g(\phi)$ .

## Separation of Variables

- ▶ The partial differential equations of  $R(r)$ ,  $f(\theta)$ , and  $g(\phi)$  are

$$\frac{\partial \psi}{\partial r} = fg \frac{\partial R}{\partial r} \quad \frac{\partial \psi}{\partial \theta} = Rg \frac{\partial f}{\partial \theta} \quad \frac{\partial^2 \psi}{\partial \phi^2} = Rf \frac{\partial^2 g}{\partial \phi^2}$$

- ▶ Substitute them into main equation

$$-\frac{\sin^2 \theta}{R} \frac{\partial}{\partial r} \left( r^2 \frac{\partial R}{\partial r} \right) - \frac{2\mu}{\hbar^2} r^2 \sin^2 \theta (E - V) - \frac{\sin \theta}{f} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial f}{\partial \theta} \right) = \frac{1}{g} \frac{\partial^2 g}{\partial \phi^2}$$

- ▶ Multiply both sides by  $r^2 \sin^2 \theta / Rfg$

$$\frac{fg}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial R}{\partial r} \right) + \frac{Rg}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{Rf}{r^2 \sin^2 \theta} \frac{\partial^2 g}{\partial \phi^2} + \frac{2\mu}{\hbar^2} (E - V) Rfg = 0$$

## Solution of the Schrödinger Equation

- ▶ Only  $r$  and  $\theta$  appear on the left side and only  $\phi$  appears on the right side of Equation.
- ▶ The R.H.S. of equation is function of  $\phi$ .
- ▶ The L.H.S. is function of  $r$  and  $\theta$ .
- ▶ This implies ONLY ONE THING
- ▶ Each side needs to be equal to a constant for the equation to be true.

Set the constant  $-m_\ell^2$  equal to the right side of Equation

$$\frac{d^2 g}{d\phi^2} = -m_\ell^2 g \quad \text{----- azimuthal equation}$$

- ▶ It is convenient to choose a solution to be

$$e^{im_\ell\phi}$$

## Solution of the Schrödinger Equation

- ▶  $e^{im_\ell\phi}$  satisfies equation for any value of  $m_\ell$ .
- ▶ The solution be single valued in order to have a valid solution for any  $\phi$ , which is

$$g(\phi) = g(\phi + 2\pi) \longrightarrow e^0 = e^{2\pi im_\ell}$$
$$g(\phi = 0) = g(\phi = 2\pi)$$

- ▶  $m_\ell$  to be zero or an integer (positive or negative) for this to be true.
- ▶ Set the left side of eqn equal to  $-m_\ell^2$  and rearrange it.

$$\frac{1}{R} \frac{\partial}{\partial r} \left( r^2 \frac{\partial R}{\partial r} \right) + \frac{2\mu r^2}{\hbar^2} (E - V) = \frac{m_\ell^2}{\sin^2 \theta} - \frac{1}{f \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial f}{\partial \theta} \right)$$

- ▶ Everything depends on  $r$  on the left side and  $\theta$  on the right side of the equation.

## Solution of the Schrödinger Equation

- ▶ Set each side equal to constant  $\ell(\ell + 1)$ .

$$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{2\mu}{\hbar^2} \left[ E - V - \frac{\hbar^2 \ell(\ell + 1)}{2\mu r^2} \right] R = 0$$

-----Radial equation

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{df}{d\theta} \right) + \left[ \ell(\ell + 1) - \frac{m_\ell^2}{\sin^2 \theta} \right] f = 0$$

-----Angular equation

$$\frac{d^2 g}{d\phi^2} = -m_\ell^2 g$$

-----Angular equation

- ▶ Schrödinger equation has been separated into three ordinary second-order differential equations each containing only one variable.

## Solution of the Radial Equation

- ▶ The radial equation is called the associated Laguerre equation and the *solutions*  $R$  that satisfy the appropriate boundary conditions are called *associated Laguerre functions*.
- ▶ Assume the ground state has  $\ell = 0$  and this requires  $m_\ell = 0$ .

$$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{2\mu}{\hbar^2} (E - V) R = 0$$

- ▶ The derivative of  $r^2 \frac{dR}{dr}$  yields two terms.  
Write those terms and insert in equation

$$\frac{d^2 R}{dr^2} + \frac{2}{r} \frac{dR}{dr} + \frac{2\mu}{\hbar^2} \left( E + \frac{e^2}{4\pi\epsilon_0 r} \right) R = 0$$



## Solution of the Radial Equation

- ▶ Try a solution  $R = Ae^{-r/a_0}$   
 $A$  is a normalized constant.  
 $a_0$  is a constant with the dimension of length.  
Take derivatives of  $R$  and insert them into Eqn

$$\left( \frac{1}{a_0^2} + \frac{2\mu}{\hbar^2} E \right) + \left( \frac{2\mu e^2}{4\pi\epsilon_0 \hbar^2} - \frac{2}{a_0} \right) \frac{1}{r} = 0$$

- ▶ To satisfy this Eqn for any  $r$  is for each of the two expressions in parentheses to be zero.  
Set the second parentheses equal to zero and solve for  $a_0$ .

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{\mu e^2}$$

Set the first parentheses equal to zero and solve for  $E$ .

$$E = -\frac{\hbar^2}{2\mu a_0^2} = -E_0$$

Both equal to the Bohr result.

## Quantum Numbers

- The appropriate boundary conditions to and (7.11) leads to the following restrictions on the quantum numbers  $\ell$  and  $m_\ell$ :
  - $\ell = 0, 1, 2, 3, \dots$
  - $m_\ell = -\ell, -\ell + 1, \dots, -2, -1, 0, 1, 2, \dots, \ell - 1, \ell$
  - $|m_\ell| \leq \ell$  and  $\ell < \infty$ .
- The predicted energy level is

$$E_n = -\frac{E_0}{n^2}$$

# Hydrogen Atom Radial Wave Functions

- ▶ First few radial wave functions  $R_{n\ell}$

**Table 7.1 Hydrogen Atom Radial Wave Functions**

$n$	$\ell$	$R_{n\ell}(r)$
1	0	$\frac{2}{(a_0)^{3/2}} e^{-r/a_0}$
2	0	$\left(2 - \frac{r}{a_0}\right) \frac{e^{-r/2a_0}}{(2a_0)^{3/2}}$
2	1	$\frac{r}{a_0} \frac{e^{-r/2a_0}}{\sqrt{3}(2a_0)^{3/2}}$
3	0	$\frac{1}{(a_0)^{3/2}} \frac{2}{81\sqrt{3}} \left(27 - 18\frac{r}{a_0} + 2\frac{r^2}{a_0^2}\right) e^{-r/3a_0}$
3	1	$\frac{1}{(a_0)^{3/2}} \frac{4}{81\sqrt{6}} \left(6 - \frac{r}{a_0}\right) \frac{r}{a_0} e^{-r/3a_0}$
3	2	$\frac{1}{(a_0)^{3/2}} \frac{4}{81\sqrt{30}} \frac{r^2}{a_0^2} e^{-r/3a_0}$

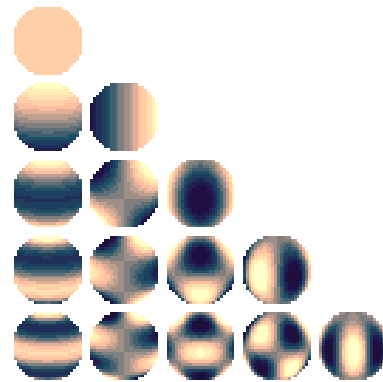
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- ▶ Subscripts on  $R$  specify the values of  $n$  and  $\ell$ .

# Solution of the Angular and Azimuthal Equations

- ▶ The solutions for Azimuthal equation are  $e^{im_\ell\phi}$  or  $e^{-im_\ell\phi}$ .
- ▶ Solutions to the angular and azimuthal equations are linked because both have  $m_\ell$ .
- ▶ Group these solutions together into functions.

$$Y(\theta, \phi) = f(\theta)g(\phi) \quad \text{----- spherical harmonics}$$



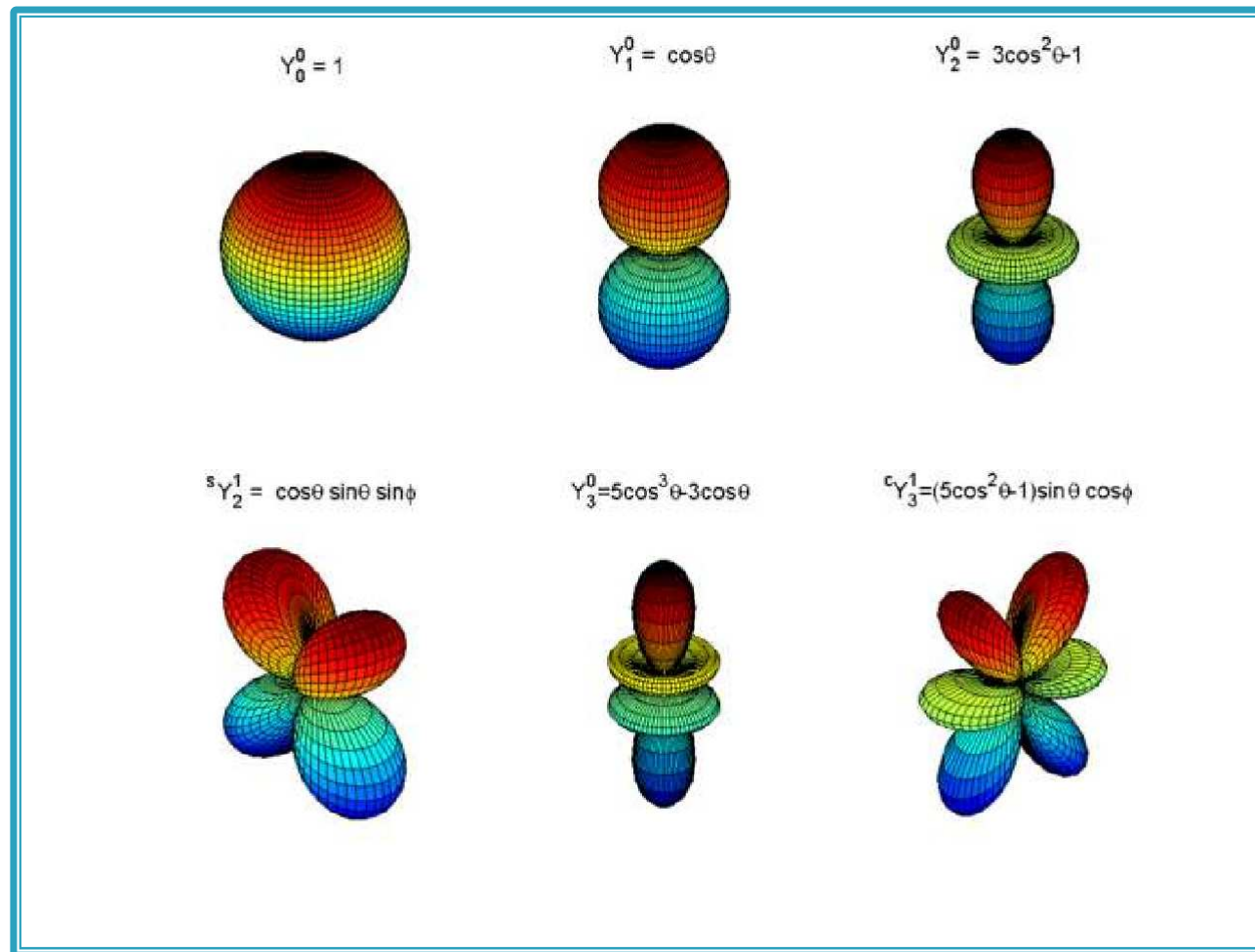
# Normalized Spherical Harmonics

**Table 7.2** Normalized Spherical Harmonics  $Y_{\ell m_{\ell}}(\theta, \phi)$

$\ell$	$m_{\ell}$	$Y_{\ell m_{\ell}}$
0	0	$\frac{1}{2\sqrt{\pi}}$
1	0	$\frac{1}{2}\sqrt{\frac{3}{\pi}} \cos \theta$
1	$\pm 1$	$\mp \frac{1}{2}\sqrt{\frac{3}{2\pi}} \sin \theta e^{\pm i\phi}$
2	0	$\frac{1}{4}\sqrt{\frac{5}{\pi}} (3 \cos^2 \theta - 1)$
2	$\pm 1$	$\mp \frac{1}{2}\sqrt{\frac{15}{2\pi}} \sin \theta \cos \theta e^{\pm i\phi}$
2	$\pm 2$	$\frac{1}{4}\sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{\pm 2i\phi}$
3	0	$\frac{1}{4}\sqrt{\frac{7}{\pi}} (5 \cos^3 \theta - 3 \cos \theta)$
3	$\pm 1$	$\mp \frac{1}{8}\sqrt{\frac{21}{\pi}} \sin \theta (5 \cos^2 \theta - 1) e^{\pm i\phi}$
3	$\pm 2$	$\frac{1}{4}\sqrt{\frac{105}{2\pi}} \sin^2 \theta \cos \theta e^{\pm 2i\phi}$
3	$\pm 3$	$\mp \frac{1}{8}\sqrt{\frac{35}{\pi}} \sin^3 \theta e^{\pm 3i\phi}$

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# Normalized Spherical Harmonics



## Solution of the Angular and Azimuthal Equations

- ▶ The radial wave function  $R$  and the spherical harmonics  $Y$  determine the probability density for the various quantum states. The total wave function  $\psi(r, \theta, \phi)$  depends on  $n$ ,  $\ell$ , and  $m_\ell$ . The wave function becomes

$$\Psi_{nlm_\ell}(r, \theta, \phi) = R_{nl}(r)Y_{lm_\ell}(\theta, \phi)$$

# Quantum Numbers

The three quantum numbers:

- $n$  Principal quantum number
- $\ell$  Orbital angular momentum quantum number
- $m_\ell$  Magnetic quantum number

The boundary conditions:

- $n = 1, 2, 3, 4, \dots$  Integer
- $\ell = 0, 1, 2, 3, \dots, n - 1$  Integer
- $m_\ell = -\ell, -\ell + 1, \dots, 0, 1, \dots, \ell - 1, \ell$  Integer

The restrictions for quantum numbers:

- $n > 0$
- $\ell < n$
- $|m_\ell| \leq \ell$



## Principal Quantum Number $n$

- ▶ It results from the solution of  $R(r)$  because  $R(r)$  includes the potential energy  $V(r)$ .

The result for this quantized energy is

$$E_n = \frac{-\mu}{2} \left( \frac{e^2}{4\pi\epsilon_0\hbar} \right)^2 \frac{1}{n^2} = -\frac{E_0}{n^2}$$

- ▶ The negative means the energy  $E$  indicates that the electron and proton are bound together.

## Orbital Angular Momentum Quantum Number $\ell$

- ▶ It is associated with the  $R(r)$  and  $f(\theta)$  parts of the wave function.
- ▶ Classically, the orbital angular momentum  $\vec{L} = \vec{r} \times \vec{p}$  with  $L = mv_{\text{orbital}}r$ .
- ▶  $\ell$  is related to  $L$  by  $L = \sqrt{\ell(\ell+1)}\hbar$  .
- ▶ In an  $\ell = 0$  state,  $L = \sqrt{0(1)}\hbar = 0$

It disagrees with Bohr's semiclassical "planetary" model of electrons orbiting a nucleus  $L = n\hbar$ .

# Orbital Angular Momentum Quantum Number $\ell$

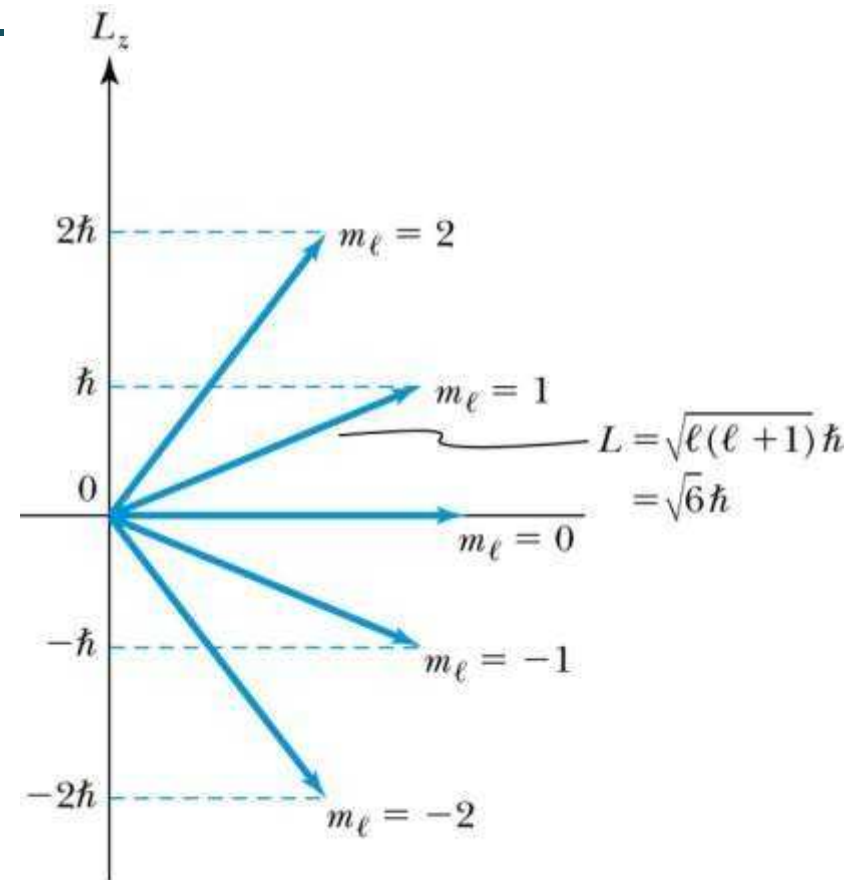
- ▶ A certain energy level is degenerate with respect to  $\ell$  when the energy is independent of  $\ell$ .
- ▶ Use letter names for the various  $\ell$  values.
  - $\ell =$         0        1        2        3        4        5 ...
  - Letter =    *s*        *p*        *d*        *f*        *g*        *h*...
- ▶ Atomic states are referred to by their  $n$  and  $\ell$ .
- ▶ A state with  $n = 2$  and  $\ell = 1$  is called a  $2p$  state.
- ▶ The boundary conditions require  $n > \ell$ .

## Magnetic Quantum Number $m_\ell$

- The angle  $\Phi$  is a measure of the rotation about the z axis.
- The solution for  $g(\Phi)$  specifies that  $m_\ell$  is an integer and related to the z component of  $L$ .

$$L_z = m_\ell \hbar$$

- The relationship of  $L$ ,  $L_z$ ,  $\ell$ , and  $m_\ell$  for  $\ell = 2$ .
- $L = \sqrt{\ell(\ell+1)}\hbar = \sqrt{6}\hbar$  is fixed because  $L_z$  is quantized.
- Only certain orientations  $\vec{L}$  are possible and this is called space quantization.



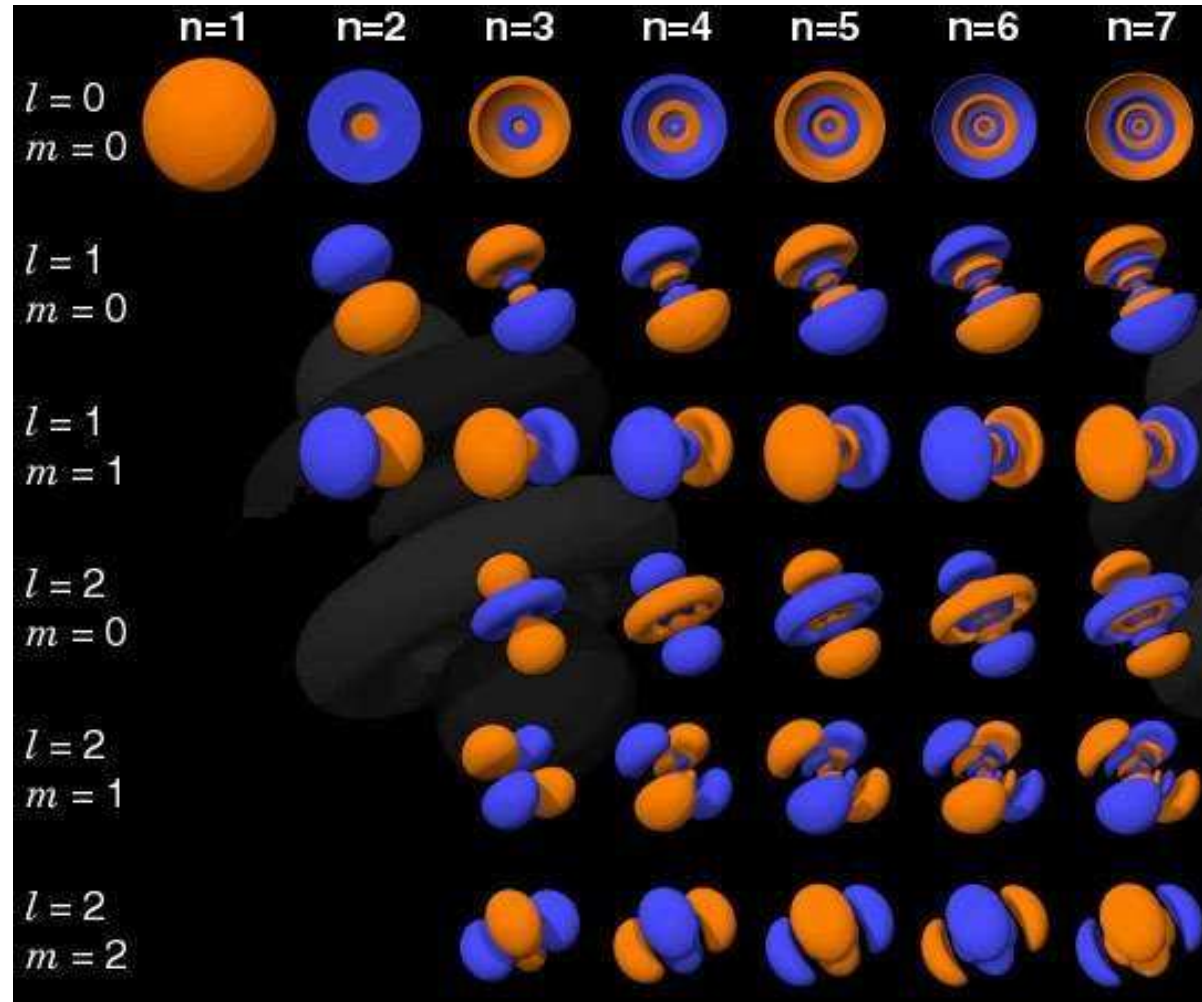
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## Magnetic Quantum Number $m_\ell$

- Quantum mechanics allows  $\vec{L}$  to be quantized along only one direction in space. Because of the relation  $L^2 = L_x^2 + L_y^2 + L_z^2$  the knowledge of a second component would imply a knowledge of the third component because we know  $\vec{L}$ .
- We expect the average of the angular momentum components squared to be  $\langle L_x^2 \rangle = \langle L_y^2 \rangle = \langle L_z^2 \rangle$ .

$$\langle L^2 \rangle = 3 \langle L_z^2 \rangle = \frac{3}{2\ell + 1} \sum_{m_\ell = -\ell}^{\ell} m_\ell^2 \hbar^2 = \ell(\ell + 1) \hbar^2$$

# Hydrogen Orbitals



**Thank You!!!**

