

VISCOUS FLOW OF INCOMPRESSIBLE FLOW

(LAMINAR FLOW)

• Reynold's Number (Re) :-

It is defined as the ratio of Inertia force to viscous force.

$$Re = \frac{\text{Inertia force}}{\text{Viscous force}} = \frac{F_i}{F_v} = \frac{\rho V L}{\mu} = \frac{VL}{\mu/\rho}$$

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|-----------------------|
| $Re = \frac{VL}{\nu}$ |
|-----------------------|

For pipe flow,

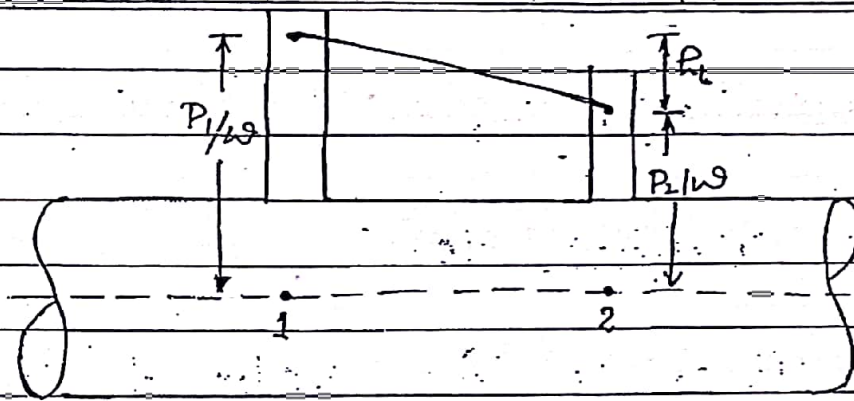
$Re < 2000 \rightarrow$ Laminar (or viscous) flow.

$Re > 4000 \rightarrow$ Turbulent flow.

$2000 < Re < 4000 \rightarrow$ Transition flow.

NOTE- For pipe flow the characteristic length L is taken as dia. of the pipe.

| |
|-----------------------|
| $Re = \frac{VD}{\nu}$ |
|-----------------------|



Apply continuity eqn. b/w 1 & 2.

$$Q = A_1 V_1 = A_2 V_2$$

$$V_1 = V_2$$

$$[\because A_1 = A_2]$$

Apply Bernoulli eqn. b/w 1 & 2.

$$\frac{P_1}{w} + \frac{v_1^2}{2g} = \frac{P_2}{w} + \frac{v_2^2}{2g} + h_L$$

$$\frac{P_1}{w} = \frac{P_2}{w} + h_L$$

$$h_L = \frac{P_1 - P_2}{w}$$

In the direction of flow press. decreases in overcome to losses because press. gradient in the direction of flow is negative.

Darcy-Weisbach Equation

This equation is used for finding out head loss due to friction in steady laminar & turbulent flow.

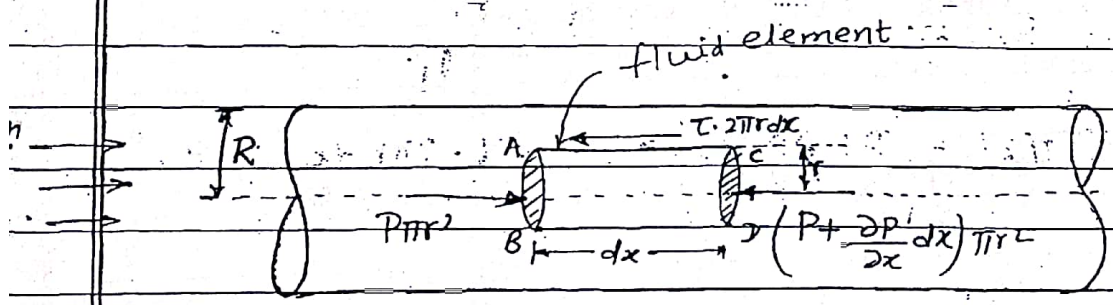
$$h_f = \frac{4f'LV^2}{2gD} \quad \leftarrow \text{oldest formula.}$$

where f' = friction co-efficient.

| | |
|---------------------------|--------------------------------------|
| $h_f = \frac{fLV^2}{2gD}$ | \leftarrow new formula. (use this) |
|---------------------------|--------------------------------------|

where f = friction factor = $4f'$

Laminar flow through circular pipe (Hagen-Poiseuille flow)



Consider a horizontal pipe of radius R . The viscous fluid is flowing from left to right in the pipe as shown in fig.

Consider a fluid element of radius r , sliding in

A cylindrical fluid element of radius $(r+dr)$:
let the length of the fluid element be dx .

If 'P' is the intensity of pressure on the face AB, then intensity of pressure on face CD will be $(P + \frac{\partial P}{\partial x} dx)$. [from Euler's Expansion series]

Then the force acting on fluid element -

1. The press. force, $P \times \pi r^2$ on face AB.
2. The press. force, $(P + \frac{\partial P}{\partial x} dx) \pi r^2$ on face CD.
3. The shear force, $\tau \times 2\pi r dx$ on the surface of fluid element.

As there is no acceleration, hence the summation of all forces in the direction of flow must be zero i.e.

$$P \pi r^2 - \left(P + \frac{\partial P}{\partial x} dx \right) \pi r^2 - \tau \cdot 2\pi r dx = 0$$

$$Pr - \left(P + \frac{\partial P}{\partial x} dx \right) r = 2\tau \cdot dx$$

$$Pr - Pr - \frac{\partial P}{\partial x} dx \cdot r = 2\tau \cdot dx$$

$$\frac{-\partial P}{\partial x} dx \cdot r = 2\tau dx$$

$$\tau = \left(\frac{-\partial P}{\partial x} \right) \frac{r}{2}$$

NOTE- The shear stress τ across a section varies with 'r' as $\frac{\partial P}{\partial x}$ across the section is constant. Hence shear stress distribution across a section is linear.

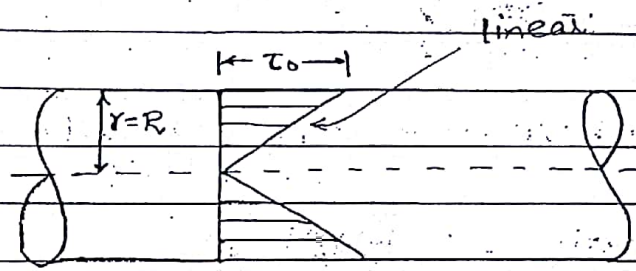
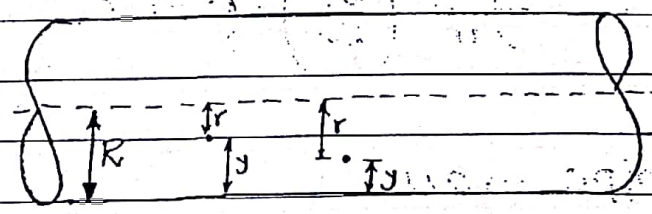


fig. shear stress distribution across a section.

maximum shear stress τ_0 at the pipe wall.

velocity distribution



$\therefore R = r + y$

[$\because R$ is const. differentiate the eqn. we get]

$0 = dr + dy$

$\Rightarrow dy = -dr$

but we know that,

$$\tau = \mu \cdot \frac{du}{dy}$$

$$\tau = \mu \frac{du}{(-dr)}$$

$$\tau = -\mu \frac{du}{dr}$$

& also $\tau = \left(\frac{-\partial P}{\partial x} \right) \frac{r}{2}$

Equating above two eqn, we get:

$$-\mu \frac{du}{dr} = \left(\frac{-\partial P}{\partial x} \right) \frac{r}{2}$$

$$du = \frac{1}{2\mu} \left(\frac{\partial P}{\partial x} \right) \cdot r \cdot dr$$

Integrate the above eqn. we find:

$$u = \frac{1}{2\mu} \left(\frac{\partial P}{\partial x} \right) \frac{r^2}{2} + C$$

At the pipe wall,

$$r = R \quad \& \quad u = 0$$

then, $0 = \frac{1}{2\mu} \left(\frac{\partial P}{\partial x} \right) \frac{R^2}{2} + C$

$$C = \frac{-1}{4\mu} \left(\frac{\partial P}{\partial x} \right) R^2$$

then, $u = \frac{1}{4\mu} \left(\frac{\partial P}{\partial x} \right) \frac{r^2}{2} = \frac{1}{4\mu} \left(\frac{\partial P}{\partial x} \right) R^2$

$$u = \frac{-1}{4\mu} \left(\frac{\partial P}{\partial x} \right) R^2 \left[\frac{1-r^2}{R^2} \right]$$

At the centre,

$$r=0 \quad \& \quad u = u_{max}$$

$$u_{max} = \frac{-1}{4\mu} \left(\frac{\partial P}{\partial x} \right) R^2 \quad \text{*****}$$

& $u = u_{max} \left(\frac{1-r^2}{R^2} \right) \quad \text{*****}$

NOTE- value of μ , $\frac{\partial P}{\partial x}$ & R are constant, which means the velocity, u varies with the square of r . Thus the eqn. is parabolic eqn. This shows the velocity distribution across the section of a pipe is parabolic.

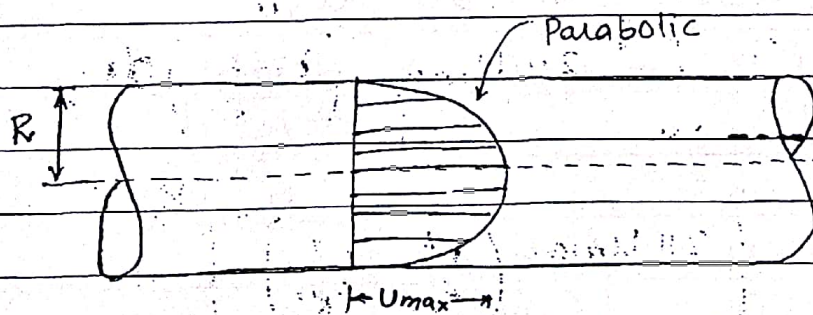
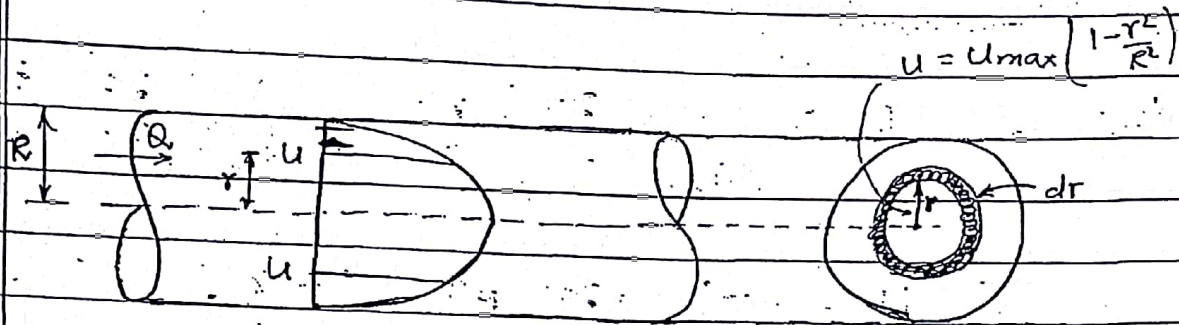


fig- velocity distribution across a section

• Discharge (Q) :-



$$u = u_{max} \left(1 - \frac{r^2}{R^2}\right) \quad \therefore u \text{ vary with } r$$

local velocity

discharge through circular ring element of radius r & thickness dr is,

$dQ = \text{velocity at radius } r \times \text{area of ring element}$

$$dQ = u \times 2\pi r dr$$

$$= u_{max} \left(1 - \frac{r^2}{R^2}\right) \times 2\pi r dr$$

Integrate the above eqn. we get,

$$\text{then, } Q = \int_0^R 2\pi u_{max} \left(\frac{r - r^3}{R^2}\right) dr.$$

$$Q = \left[2\pi u_{max} \left(\frac{r^2}{2} - \frac{r^4}{4R^2}\right) \right]_0^R$$

$$Q = 2\pi U_{max} \left[\frac{R^2}{2} - \frac{R^4}{4R^2} \right]$$

$$Q = 2\pi U_{max} \left[\frac{R^2}{4} \right]$$

$$Q = \frac{\pi U_{max} R^2}{2}$$

$$Q = \frac{\pi R^2}{2} \times \frac{1}{4\mu} \left(\frac{-\partial P}{\partial x} \right) R^2$$

$$Q = \frac{\pi}{8\mu} \left(\frac{-\partial P}{\partial x} \right) R^4$$

• Average velocity :-

$$Q = AV$$

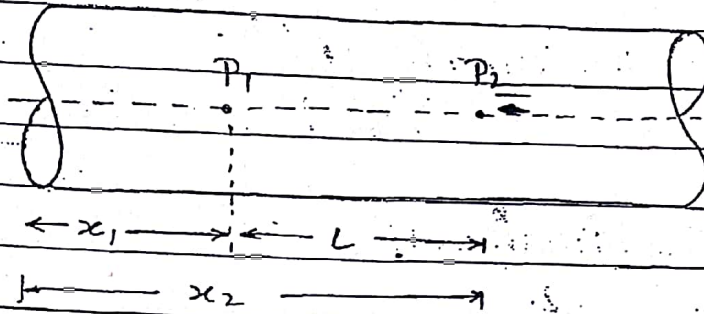
$$V = \frac{Q}{A}$$

$$V = \frac{\pi U_{max} R^2}{\pi R^2}$$

$$V = \frac{U_{max}}{2}$$

It is defined as the ratio of total discharge to the total area of flow.

Pressure drop in a given length (L) :-



$$V = \frac{U_{\max}}{2}$$

$$V = \frac{1}{2} \left[\frac{1}{4\mu} \left(\frac{-\partial P}{\partial x} \right) R^2 \right]$$

$$\frac{8\mu V}{R^2} \frac{\partial x}{\partial x} = -\partial P$$

Integrate the above eqn.

$$\int_{x_1}^{x_2} \frac{8\mu V}{R^2} \partial x = \int_{P_1}^{P_2} -\partial P$$

$$\frac{8\mu V}{R^2} (x_2 - x_1) = -(P_2 - P_1)$$

$$\frac{8\mu V L}{R^2} = P_1 - P_2$$

$$P_1 - P_2 = \frac{8\mu V L}{R^2}$$

$$P_1 - P_2 = \frac{8 \mu V L}{(D/2)^2}$$

$$P_1 - P_2 = \frac{32 \mu V L}{D^2}$$

avg. velocity *****

Heart of the laminar flow eqn.

We know that,

$$\frac{P_1 - P_2}{L} = \frac{4 \mu v}{r^2}$$

$$P_1 - P_2 = 4 \mu v L$$

$$P_1 - P_2 = \rho g h r$$

$$\frac{32 \mu V L}{D^2} = \rho g \times \frac{f L V^2}{2 g D}$$

$$\frac{32 \mu}{D} = \frac{\rho f V^2}{2}$$

$$\frac{64 \mu}{\rho V D} = f$$

$$f = \frac{64 \mu}{\rho V D}$$

$$f = \frac{64}{Re}$$

Imp.

NOTE- Friction factor in laminar flow through circular pipe depends upon Reynold number only.

∴ shear velocity (v^*) :-

We know that,

$$\tau = \left(\frac{-\partial P}{\partial x} \right) \frac{r}{2}$$

At the wall of the pipe,

$$\tau = \tau_0 \quad \text{at } r = R$$

then,
$$\tau_0 = \left(\frac{-\partial P}{\partial x} \right) \frac{R}{2}$$

$$\tau_0 = - \frac{(P_2 - P_1)}{x_2 - x_1} \cdot \frac{R}{2}$$

$$\tau_0 = \frac{P_1 - P_2}{x_2 - x_1} \cdot \frac{D}{4}$$

$$\tau_0 = \frac{P_1 - P_2}{L} \cdot \frac{D}{4} \quad [\because x_2 - x_1 = L]$$

$$\tau_0 = (P_1 - P_2) \cdot \frac{D}{4L}$$

But we know that,

$$\frac{P_1 - P_2}{w} = h_L$$

$$\Rightarrow P_1 - P_2 = w h_L$$

then, $T_0 = w h_L \cdot \frac{D}{4L}$

$$T_0 = \frac{\rho g \times f L V^2 \cdot D}{2gD \cdot 4L}$$

$$T_0 = \frac{\rho f V^2}{8}$$

$$\Rightarrow \frac{T_0}{\rho} = \frac{f}{8} V^2$$

$$\sqrt{\frac{T_0}{\rho}} = \sqrt{\frac{f}{8}} V$$

$$V^* = \sqrt{\frac{f}{8}} V$$

Where $\sqrt{\frac{T_0}{\rho}} = V^* = \text{shear velocity.}$